



Effect of Electrode and Electrolyte Thicknesses on All-Solid-State Battery Performance Analyzed With the Sand Equation

Didier Devaux^{1*}, Hugo Leduc¹, Philippe Dumaz¹, Margaux Lecuyer², Marc Deschamps² and Renaud Bouchet^{1*}

¹ Univ. Grenoble Alpes, Univ. Savoie Mont Blanc, CNRS, Grenoble INP, LEPMI, Grenoble, France, ² Blue Solutions, Odet, Ergué Gabéric, Quimper, France

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*Correspondence:

Didier Devaux
didier.devaux@grenoble-inp.fr
Renaud Bouchet
renaud.bouchet@grenoble-inp.fr

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The energy conversion and storage are great challenges for our society. Despite the progress accomplished by the Lithium(Li)-ion technology based on flammable liquid electrolyte, their intrinsic instability is the strong safety issue for large scale applications. The use of solid polymer electrolytes (SPEs) is an adequate solution in terms of safety and energy density. To increase the energy density (resp. specific energy) of the batteries, the positive electrode thickness must be augmented. However, as for Li-ion liquid electrolyte, the cationic transference number of SPEs is low, typically below 0.2, which limits their power performance because of the formation of strong gradient of concentration throughout the battery. Thus, for a given battery system a compromise between the energy density and the power has to be found in a rapid manner. The goal of this study is to propose a simple efficient methodology to optimize the thickness of the SPE and the positive electrode based on charge transport parameters, which allows to determine the effective limiting Li⁺ diffusion coefficient. First, we rapidly establish the battery power performance thanks to a specific discharge protocol. Then, by using an approach based on the Sand equation a limiting current density is determined. A unique mother curve of the capacity as a function of the limiting current density is obtained whatever the electrode and electrolyte thicknesses. Finally, the effective limiting diffusion coefficient is estimated which in turn allows to design the best electrode depending on electrolyte thickness.

Keywords: lithium battery, polymer electrolyte, power performance, transport limitation, effective diffusion coefficient, power cycling procedure

INTRODUCTION

Batteries are one of the most widely used electrochemical energy storage devices thanks to their high energy permitting to operate devices for a long period of time (Kim et al., 2015). The Lithium (Li)-ion system is the predominant solution to power many applications from small electronic devices to electric cars (Kurzweil and Garche, 2017). As the demand in energy output from consumers is constantly increasing new battery systems must be developed and optimized depending on the application requirements. In this context Li metal is an ideal as negative electrode due to its high specific capacity, and low operating voltage (Xu et al., 2014). However, in combination with a conventional liquid electrolyte, the Li electrodeposits on the metallic Li are uneven giving rises

to low Coulombic efficiency and dendrite growth (Tarascon and Armand, 2001; Wu et al., 2019). Replacing the liquid electrolyte by a solid polymer electrolyte (SPE) permits to envision safe high energy density batteries (Armand, 1994; Agrawal and Pandey, 2008).

In term of design, a practical Li metal battery stack should be optimized through its positive electrode formulation and overall material assembly notably to reduce the cost per elemental cell (Gallagher et al., 2014). The goal being to produce a long lasting battery with high rate performance to obtain extended specific energy (Eftekhari, 2017; Schnell et al., 2018; Zeng et al., 2019). Regarding the SPE, enough ionic conductivity in the order of 1 mS.cm^{-1} is obtained at 80°C for PEO-based electrolyte (Devaux et al., 2012). PEO possesses mechanical properties and flexibility high enough to be processed by hot-pressing, extrusion, or solvent-casting methods to form of thin films ranging from 10 to $100 \mu\text{m}$ (Baudry et al., 1997; Porcarelli et al., 2016; Schnell et al., 2018; Yang et al., 2019). The goals being to minimize the SPE thickness to reduce ohmic loss and to increase the diffusion limited current density during battery operation. Moreover, PEO possesses other advantages such as a good adhesion, and electrochemical and chemical stability toward metallic Li (Armand, 1983). Indeed, PEO is usually used as a fundamental SPE brick which will also contain other polymers or additives to improve the desired properties-transport number, mechanical strength, electrochemical window, etc. (Bouchet et al., 2013; Zhang et al., 2017). LiFePO_4 is a reference candidate as positive active materials thanks to its bi-phased redox reaction at about $3.43 \text{ V vs. Li}^+/\text{Li}^\circ$ which ensures safety (Padhi et al., 1997). At last, for a given electrode formulation, the energy density is directly linked to the active material loading (Du et al., 2017). So, for a given active material, increasing the electrode thickness is the simplest way to reach higher specific energy.

The electrode formulation is of importance through the loading of active material and the electrode porosity to reach the best power performance at an optimal electrode thickness (Newman, 1995; Yu et al., 2006, 2013). Zheng et al. reported on positive electrodes made of $\text{LiNi}_{1/3}\text{Co}_{1/3}\text{Mn}_{1/3}\text{O}_2$ (NCM) or LiFePO_4 and showed that Li ion diffusion within the electrolyte in the electrode was the limiting discharge process (Zheng et al., 2012). Interestingly, they showed that the specific C-rate at which capacity starts to drop is representative to the overall battery rate performance. The electrode capacity, and thus active material loading, follows a negative power law, with the C-rate. Therefore, as expected low loaded electrodes provide higher rate performance than their higher loaded counterparts. A similar conclusion was substantially drawn by Gallagher et al. who combined experimental and simulation approaches and Heubner et al. on thick NCM electrodes (Gallagher et al., 2016; Heubner et al., 2019). In consequence, the ionic diffusion processes at stake in a battery system must be fully characterized as it governs the battery voltage divergence before the full recovery of the capacity because the Li-ion concentration reaches a null concentration at the cathode. All other electrochemical processes (ion and electron migration, charge transfer, etc.) are mostly related to “ohmic” drops which simply increase the battery polarization when the rate increases.

Typically, the relationship between discharge capacity and C-rate is similar for each battery technology. At low C-rate, the capacity is maximal and constant. At C-rate higher than a threshold value, the capacity falls rapidly (Gallagher et al., 2016). This behavior limits the high rate capabilities in terms of rapid charge and fast discharge (Doyle and Newman, 1997; Kang and Ceder, 2009). In literature a series of empirical and semi-empirical equations based on a phenomenological stretched exponential decay function have been developed to fit the sigmoidal shape of the capacity-C/rate curve (Gallagher et al., 2016; Wong et al., 2017; Heubner et al., 2018b; Tian et al., 2019). The physical meaning behind these equations is to consider a time dependent relaxation process describing herein a single rate limiting diffusion phenomena. So, those equations use as main parameters (i) a time constant and (ii) an empirical exponent value to stretch the exponential function. The meaning of the exponent is unclear as some studies set this parameter to a constant value of 2 or let it free to get better fits. The uncertainty regarding this empirical parameter is similar to that encounter with the Peukert exponent (Doerffel and Sharkh, 2006).

Therefore, ionic diffusion in the electrolyte, and within the active materials are known to be the major physical limitations that affect the battery rate capability (Jiang and Peng, 2016). Thus by coupling experimental and simulation analysis, many studies aim to improve power and energy from the single cell to the stack level (Yu et al., 2006; Safari and Delacourt, 2011; Singh et al., 2015). The mostly accepted theoretical model is the one developed by Newman and coworker based on concentrated solution theory (Doyle et al., 1993; Fuller et al., 1994; Newman, 1995; Srinivasan and Newman, 2004). A complete description of a given battery system is required to fit the cycling data which requires numerous parameters. Some of these parameters such as transference number, ionic diffusion coefficients, thermodynamic factors or tortuosity, to cite a few, are difficult to obtain properly as their determination is rather tough and non-univocal which can lead to large approximation. In addition, each set of parameters has to be determined again each time a chemical modification is made (additives, binders, electrode formulation, active materials...) which can be time consuming.

In practice, the fast determination of the discharge capacity depending on the C-rate is of interest to screen series of batteries in which many parameters are tuned such as the porosity, the formulation, and the active material loading of electrodes or even the nature of the current collectors. Back in 1994, a protocol was proposed composed of successive discharges starting from the highest C-rate down to the lowest without charging step in between (Doyle et al., 1994). More recently, Heubner et al. presented a methodology based on chronoamperometry measurement (Heubner et al., 2018a). Starting from charged Li-ion batteries, instead of applying a constant current as discharge step a constant voltage step at the lower battery cut-off voltage is performed. The transitory current recorded is then converted through the current-time integration into relative capacity. This method, a slow charge plus the chronoamperometry, takes about a dozen of hours to complete, but presents some limitations to fully capture the battery capacity when the active materials go

through different phase change in charge (resp. in discharge). As pointed out by the authors, a typical example of such material is LiFePO_4 with a two-phase mechanism.

Herein, firstly, to reduce the time-consuming measures necessary to test battery power depending on all the parameters a cycling procedure similar to that of Doyle et al. is used (Doyle et al., 1994). Through this rapid cycling routine, the relationship between discharge capacity and current density is determined with a good accuracy at least 8 times faster than conventional galvanostatic charge-discharge cycles alternation. Then, we propose two complementary methodologies to quickly estimate the effective diffusion coefficient of the rate limiting process at stake which allows by comparing it to the diffusion coefficient into the electrolyte or active material particles, to identify which component can be optimized. Those methodologies are based on limiting current and Sand time (Sand, 1901; Bard and Faulkner, 2001). For this purpose model all-solid-state batteries made of Li metal, PEO based electrolyte acting as SPE, and LFP-based positive electrode were used. By keeping the negative electrode in excess, the interplay between positive electrode and SPE thicknesses is finely understood and general rules are provided to pinpoint toward the optimal balance in thicknesses between SPE and positive electrode. The goal being to provide a simple physical interpretation of practical power signature curves based on the Sand's equation, which permits the determination of the effective limiting diffusion processes in real system.

EXPERIMENTAL

The metallic Li foil was provided by the company Blue Solutions. The Li was stored in an Argon filled glove box (Campus, Jacomex) with sub-ppm values of H_2O and O_2 . The solid polymer electrolyte (SPE) is a random copolymer comprising PEO to ensure ionic motion and poly(propylene oxide) to provide flexibility to the resulting membrane, and doped with lithium bis-trifluoromethanesulfonimide (LiTFSI) salt. The concentration of Li salt was added to reach a molar ratio of ethylene oxide to Li salt of 25. The Li salt and Li metal foil are stored in an Argon filled glove box (Campus, Jacomex) with sub ppm value of H_2O and O_2 . The polymer and LiTFSI were dissolved in Dimethylformamide (DMF) by stirring at 350 rpm and 80°C for 3 h in a vial. When the solution was clear and transparent, it was cast onto Petri dish made of Teflon and let dry at ambient temperature (T) for 8 h. Then, to remove the solvent, the Petri dish was placed in an oven at 60°C for at least 3 days. The resulting SPE membrane was peeled off the Teflon substrate and placed inside the Ar glove box for at least a week prior any further manipulation. Still inside the glove box, SPEs were cut into wide pieces and several of them placed in a hot press. The SPEs were then pressed for several minutes à 80°C at 200 bars while adjusting the SPE quantity. After pressing and cooling, SPE disks were punch out from the pressed SPE pieces. This lead to SPE disk with thickness (y) of 18, 36, 54, 108, and $216\ \mu\text{m}$. Low porosity (<5%) positive electrodes using LiFeO_4 as active materials were specifically prepared by Blue Solutions through an extrusion process. The thickness of the positive electrode (x) were 20, 33, 48, and $60\ \mu\text{m}$ on top

of an Al current collector. The capacities of the electrodes are proportional to their thicknesses.

Inside the globe box, a Li disk of 14 mm diameter was punch out from the Li foil. Then, a layer of SPE was also cut in a disk shape of 16 mm in diameter. The two materials were laminated at 80°C and 3 bars multiple times using a home-made laminating machine until the SPE was fully adherent to the Li. The electrolyte thickness was checked after the laminating process and no variation was observed. A 12 mm diameter positive electrode disk was also punch out from the electrode foil. Several passages through the laminating machine were done to ensure adhesion between the three material layers composing the Li metal polymer battery. The battery assemblies were varied in terms of positive electrode and SPE thicknesses. So, all along the text, the Li Metal Polymer (LMP) batteries are denoted $\text{LMP}(x-y)$ with x and y corresponding to the electrode and SPE thicknesses, respectively. For each $\text{LMP}(x-y)$ system, between 4 and 6 replicates were assembled.

After assembly, the battery was placed inside a stainless steel CR2032 coin cell using stainless steel shims and a wave spring. The coin cell was sealed using a crimping machine inside the glove box. On average 1.2 bars are applied to the battery assembly with an active surface S corresponding to the positive electrode geometrical surface. Then, the battery cells were taken out of the glove box and placed on coin cell holder. The coin cells were put inside an oven (Memmert) held at 80°C and connected to a VMP3 (Bio-Logic SAS) multi-potentiostat with impedance capability.

The cycling procedure consists in a series of charge-discharge galvanostatic cycles in between 2.5 and 3.7 V vs. $\text{Li}^+/\text{Li}^\circ$. All along the text the battery potential E is referred to the $\text{Li}^+/\text{Li}^\circ$ couple. Initially, the batteries are subjected to 8 conditioning cycles at a low current density (J_0) identical in charge and discharge so that the delivered specific capacity corresponds to the effective LiFePO_4 capacity, $160\ \text{mAh}\cdot\text{g}^{-1}$. After this initial procedure, conventional galvanostatic cycling to produce the power signature of the battery is made. It consists in a series of subsequent charge and discharge steps by constantly charging at a current density of J_0 and by increasing the discharge current density J_n . In between each cycle charge/discharge a rest period of 30 min is used to relax the gradients of concentration. Taking into account the rest periods, this conventional cycling procedure needs more than a week-long time period. At each step, the areal charge capacity (Q_n) was calculated by integration of the current density J_n over time (t) during the galvanostatic steps according to:

$$Q_n = \int J_n(t) \cdot dt \quad (1)$$

For the rapid power procedure, after the conditioning galvanostatic cycles at J_0 previously described, the batteries are fully charged also at J_0 , then we accumulate several discharges separated by a rest period of 30 min to relax the gradients of concentration. We start by the highest current density, followed by 30 min of relaxation, then a discharge at slightly lower current density is performed, and so on, until a

final discharge step at J_0 (the lowest current density). This is done without charging the battery in between each discharge steps. For clarity, the cycling behavior, evolution of E as a function of the fraction δ of Li inserted in the $\text{Li}_\delta\text{FePO}_4$ phase, with $0 < \delta < 1$, for a representative LMP(48-18) battery is provided in **Supplementary Figure 1**. This rapid power procedure takes about 1 day to fully complete from the charge to the final discharge step, at least eight times shorter than the conventional cycling procedure. The underlying assumption of this cycling procedure is that by starting from a fully charged battery, the discharge capacity at a given current density J_n is the sum of the discharge capacities obtained at higher current densities, $J > J_n$, plus the one obtained when applying J_n . In other word, $Q_n(J_n)$ is calculated based on the following equation:

$$Q_n = \sum_{N \geq n} [J_N \cdot \Delta t_N] \quad (2)$$

with Δt_N the time necessary to discharge the battery under the constant current density J_N .

For both cycling procedures, the capacities calculated at each J_n of the battery replicates lie within a typical deviation below 1% which shows the very good reproducibility of our cells. This is a mandatory condition to go further into the data analysis.

Finally, to have an independent measure of the transport properties of our electrolyte, electrochemical impedance spectroscopy was executed on Li symmetric cell comprising the SPE (Bouchet et al., 2003). This technic permits to determine the different cell resistances such as the electronic (R_c), electrolyte (R_{el}), interface (R_{int}), and diffusion (R_d) resistances. Li symmetric cells were assembled via a similar lamination process than that described for the LMP batteries and sealed in CR2032 coin cells. After placing the cells in an oven at 80°C, impedance spectroscopy using an excitation signal of 40 mV in a frequency range between 10 MHz and 0.1 Hz was done. A typical Nyquist plot showing the opposite of the impedance imaginary part $[-\text{Im}(Z)]$ as a function of the real part $[\text{Re}(Z)]$ is presented in **Supplementary Figure 2**. An electrical equivalent circuit, displays in the inset of **Supplementary Figure 2**, comprising the cell resistances (R_c , R_{el} , R_{int}), a cable inductance (L_c) at high frequency, a constant phase element for the interface (CPE_{int}), and a short Warburg element (W_d) for the diffusion loop at low frequencies, permits to model the Nyquist plot. From the Warburg element two main parameters are extracted corresponding to R_d and the relaxation time (τ_r) at the maximum of the quarter of lemniscate. R_d and R_{el} are related to the cationic transference number (t^+) while τ_r is linked to the ambipolar diffusion coefficient (D_{amb}) and the SPE thickness y by the following equations (Sørensen and Jacobsen, 1982; Ross MacDonald, 1992; Bouchet et al., 2003):

$$t^+ = \frac{R_{el}}{R_{el} + R_d} \quad (3)$$

$$\tau_r = 2.54 \cdot \frac{\left(\frac{y}{2}\right)^2}{D_{amb}} \quad (4)$$

In addition, D_{amb} is linked to the Li^+ diffusion coefficient (D_{Li^+}) by the relationship.

$$D_{\text{Li}^+} = \frac{D_{amb}}{2 \cdot (1 - t^+)} \quad (5)$$

RESULTS AND DISCUSSION

The cycling behavior of a typical LMP(48-18) battery is shown in **Figure 1** that represents E as a function of the fraction δ of Li inserted in the $\text{Li}_\delta\text{FePO}_4$ phase, with $0 < \delta < 1$, recorded during conventional cycling. For clarity, only a representative charge performed at $J_0 = 0.1 \text{ mA}\cdot\text{cm}^{-2}$ is represented in **Figure 1** and some discharge current density are indicated. The battery charge presents a typical long plateau at about 3.44 V corresponding to the oxidation of LFP materials (Padhi et al., 1997) followed by a potentiostatic step at 3.7 V to reach the full electrode capacity. During discharge, a lower potential plateau at about 3.40 V, corresponding to the reduction of LFP, is observed. This plateau is less pronounced for discharge steps executed at J_n higher than $0.3 \text{ mA}\cdot\text{cm}^{-2}$ due to an increase in the battery of gradient of concentration. Moreover, these observations on the conventional cycling behavior of a typical LMP(48-18) battery remain valid for all of the other considered LMP(x - y) batteries and associated replicates.

For the conventional cycling (see **Figure 1**) and the rapid power test (see **Supplementary Figure 1**), the discharge capacities were extracted by using Equations (1) and (2), respectively. Then, each Q_n value was normalized by the discharge capacity delivered at J_0 , denoted Q_0 . Thus, **Figure 2** represents the normalized discharge capacity, ratio Q_n/Q_0 , as a function of the discharge current densities, J_n , for the two

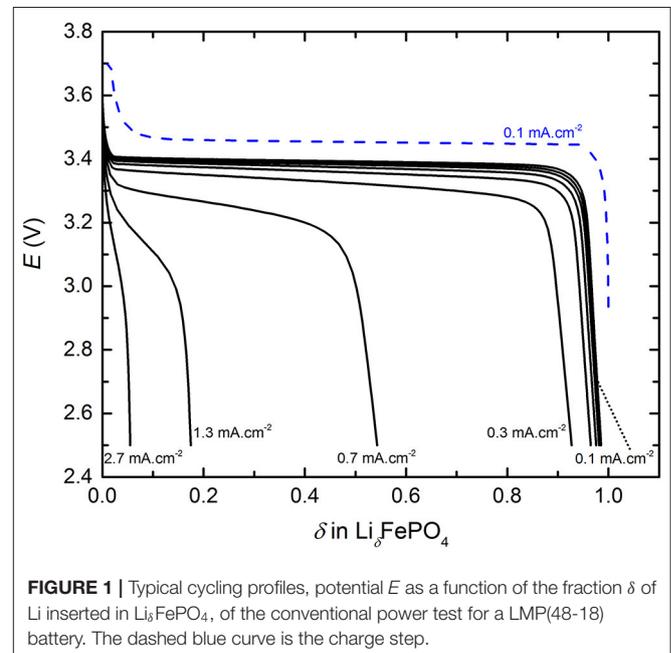
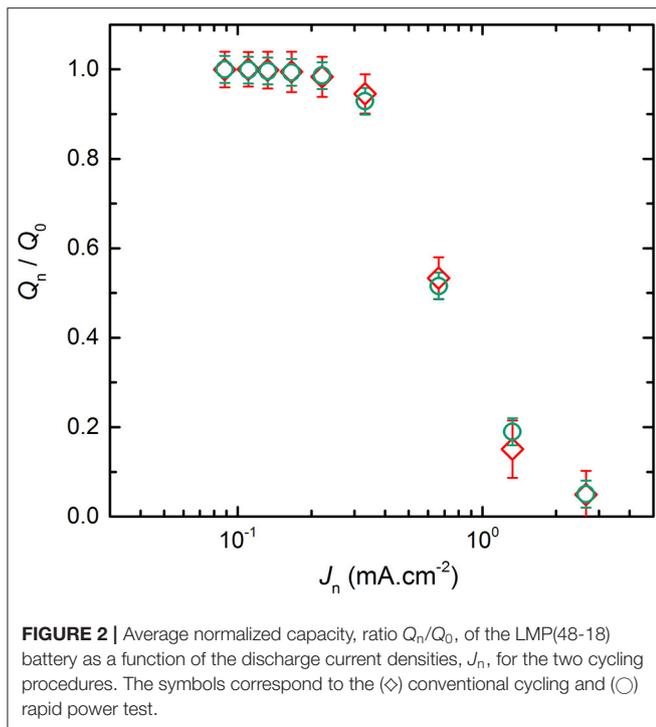
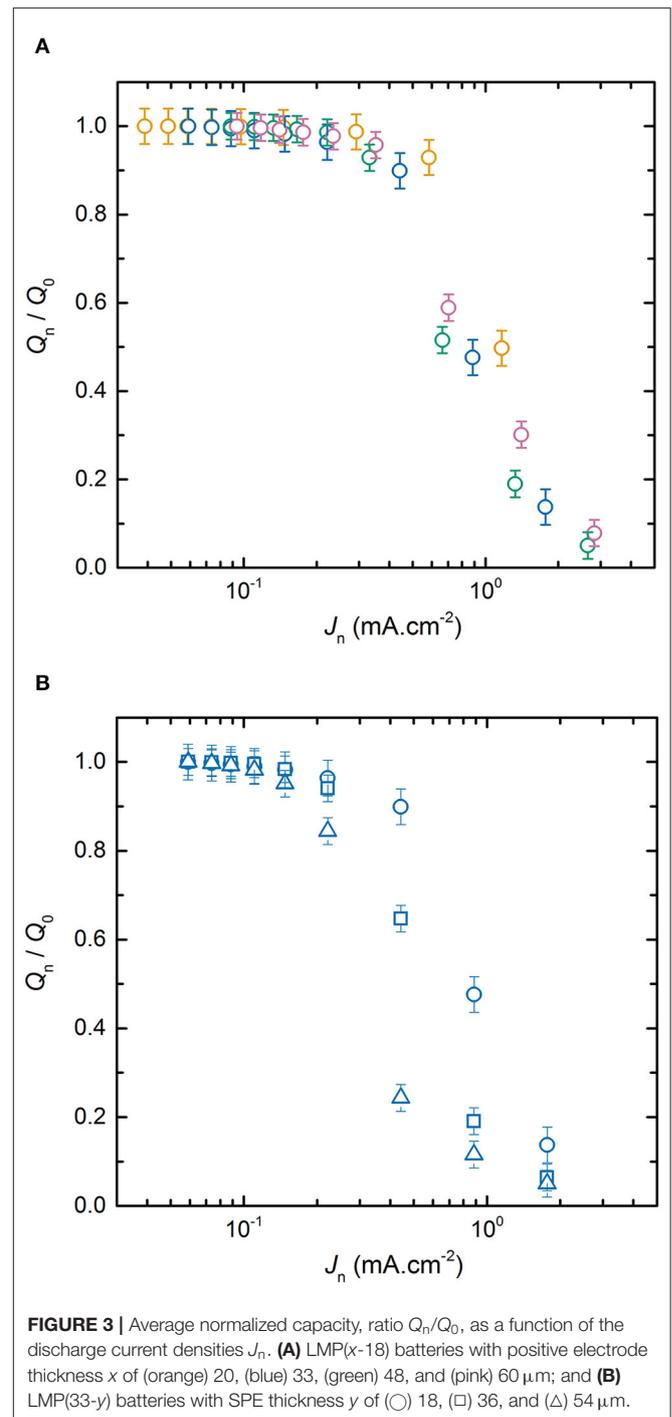


FIGURE 1 | Typical cycling profiles, potential E as a function of the fraction δ of Li inserted in $\text{Li}_\delta\text{FePO}_4$, of the conventional power test for a LMP(48-18) battery. The dashed blue curve is the charge step.



cycling procedures of the LMP(48-18) batteries. The values reported in **Figure 2** corresponds to average values with their standard deviations from the different battery replicates. For both procedures and for low J_n values, below $0.3 \text{ mA}\cdot\text{cm}^{-2}$, Q_n remains close to Q_0 by 3%. For both procedures and for high J_n values, higher than $0.3 \text{ mA}\cdot\text{cm}^{-2}$, the ratio Q_n/Q_0 quickly drops linearly with the increase of J_n down to values below 0.2 when $J_n > 1 \text{ mA}\cdot\text{cm}^{-2}$. Similarly to Li-ion batteries (Gallagher et al., 2014), the relationship between the discharge capacity and the current density of the all-solid-state Li batteries present a transition regime at a critical current density value. Moreover, the data point recorded at the highest J_n value, at $2.8 \text{ mA}\cdot\text{cm}^{-2}$, deviates from the linear trend of Q_n/Q_0 with J_n when $J_n > 0.3 \text{ mA}\cdot\text{cm}^{-2}$ due to supercapacitor effect arising from the high surface carbon fillers and coating of the LFP particles. Such capacitive effect is only seen for the highest current densities of the LMP(x - y) batteries and is not taken into account for the data analysis provided in the remaining of the text. The evolution of the normalized discharge capacities with the current densities are in excellent accordance for the conventional cycling and the rapid power test. Indeed, differences in Q_n/Q_0 values lie within the error bars typically when $J_d > 0.3 \text{ mA}\cdot\text{cm}^{-2}$. In addition, the accordance in between the two procedures is also observed for each studied LMP(x - y) batteries. The rapid power test is then a robust tool to picture the battery power performance in a time-efficient manner with a strong accuracy when compared to conventional cycling procedure. Other interesting cycling routine can be found in literature but will less accuracy regarding phase transition positive active materials as LFP (Heubner et al., 2018a).



To understand the effect of electrode thicknesses on battery power, **Figure 3A** represents the average normalized discharge capacity as a function of J_n for the LMP(x -18) batteries with x values of 20, 33, 48, and $60 \mu\text{m}$ while the SPE thickness is maintained at $18 \mu\text{m}$. For each LMP(x - y) batteries reported in **Figure 3A**, the evolution of Q_n/Q_0 with J_d is similar to that already detailed in **Figure 2**. The main difference between

each battery is the location of the transition regime between the current density regime where Q_n remains close to Q_0 and the regime where Q_n/Q_0 decreases with J_d . Based on **Figure 3A**, the thinner the positive electrode, the later Q_n will deviate strongly from Q_0 . A similar conclusion can be drawn for all the other LMP(x - y) batteries in which y is constant and x being either 20, 33, 48, or 60 μm . For completeness, **Supplementary Figure 3** shows the ratio Q_n/Q_0 as a function of J_n for the LMP(x -36) and LMP(x -54) batteries. The effect of the positive electrode thickness on the power is then similar to that reported for Li-ion batteries comprising a liquid electrolyte (Cornut et al., 2015; Heubner et al., 2018b). As a complementary case, **Figure 3B** shows the average normalized discharge capacity as a function of J_n for the LMP(33- y) batteries with y values of 18, 36, and 54 μm while the positive electrode thickness is constant at 33 μm . For a given positive electrode thickness, the transition regime occurs at lower current density when the SPE thickness increases. This observation remains valid for all the other LMP(x - y) batteries in which x is constant and y being either 18, 36, 54, 108, and 216 μm . For completeness, **Supplementary Figure 4** shows the ratio Q_n/Q_0 as a function of J_n for the LMP(20- y), LMP(48- y), and LMP(60- y). From **Figures 3A,B**, the power performance of the all-solid-state Li batteries are dependent on both positive electrode and SPE thicknesses. Interestingly, in most of the studies on standard Li-ion batteries, only the electrode thickness has been varied shading the impact of the electrolyte thickness. Understanding finely the partition of these two parameters on the battery performance is of first interest to envision an optimized battery assembly.

From each Q_n/Q_0 vs. J_n plots, we define the current density corresponding to the transition regime as the current density at which the $Q_n/Q_0 = 1$ baseline (low J_n regime) crosses the linear interpolation of Q_n/Q_0 vs. J_n (high J_n regime). This drop of capacity is due to a Li ion diffusion limitation, therefore current density corresponding to this data point is called the diffusion limiting current, denoted J_{lim} . For clarity, the graphical determination of J_{lim} is shown in **Supplementary Figure 5** for the LMP(20-18) batteries. In order to provide more insight on J_{lim} , we propose a methodology based on the Sand equation for controlled-current technique (Sand, 1901; Rosso et al., 2006). By definition, for a current density above J_{lim} , the Sand time (τ_s) corresponds to the transition time at which the flux of ionic species is not large enough to satisfy the applied current. The Sand equation (Brissot et al., 1999) is theoretically given by:

$$\tau_s = \pi \cdot D_{\text{amb}} \cdot \left(\frac{n \cdot F \cdot C_{\text{Li}}}{2 \cdot (1 - t^+) \cdot J_n} \right)^2 \quad (6)$$

with n the number of exchanged electron ($n = 1$ for LFP), F the Faraday constant ($9.648 \cdot 10^4 \text{ C} \cdot \text{mol}^{-1}$), C_{Li} the Li concentration in the SPE (Devaux et al., 2012) located in the electrolyte and in the positive electrode ($882 \text{ mol} \cdot \text{m}^{-3}$ for both), t^+ the cationic transference number, and D_{amb} the ambipolar diffusion coefficient. The use of this equation is an approximation since the boundaries condition of a plan electrode is not perfectly respected on the cathode side.

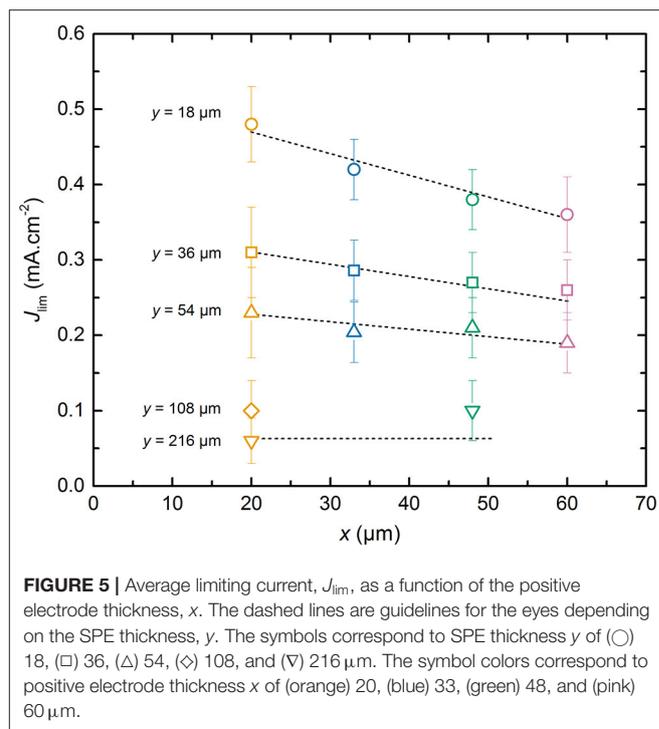
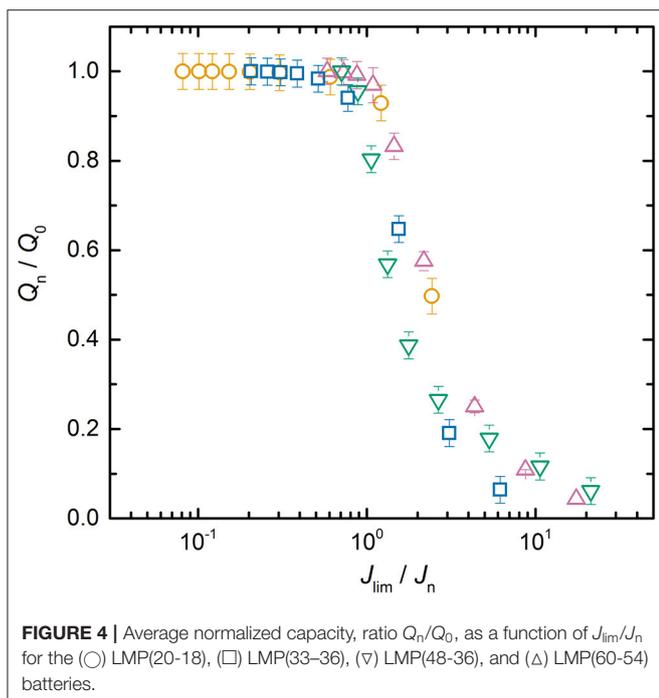
For $J_n \geq J_{\text{lim}}$, τ_s can be estimated as being equivalent to the discharge time. In other words, τ_s at a specific current density is defined as:

$$\tau_s(J_n) = \frac{Q_n}{J_n} \text{ when } J_n \geq J_{\text{lim}} \quad (7)$$

D_{amb} and t^+ of a Li salt in PEO can be measured or calculated using many techniques such as electrochemical methodologies based on polarization (Shi and Vincent, 1993; Geiculescu et al., 2006) or relaxation (Mullin et al., 2011), impedance spectroscopy (Bouchet et al., 2003), pulse-field NMR (Hayamizu et al., 1999), or molecular dynamic simulation (Diddens et al., 2010) to name a few. By looking into literature data and focusing on high molecular PEO based electrolyte, at 80°C , D_{amb} is in the range of $5 \cdot 10^{-8} \text{ cm}^2 \cdot \text{s}^{-1}$ and t^+ of about 0.15. In addition, impedance spectroscopy was performed on Li symmetric cells comprising SPE thickness of 18, 36, and 54 μm . The electrical equivalent circuit (see inset of **Supplementary Figure 2**) was used to fit all the impedance spectra in order to extract t^+ and D_{amb} values. D_{Li^+} , calculated using Equation (5), and t^+ are both independent of the SPE thickness with an average value consistent with literature analysis of $3.4 \pm 0.7 \cdot 10^{-8} \text{ cm}^2 \cdot \text{s}^{-1}$ and 0.15 ± 0.02 , respectively.

The plot of τ_s as a function of J_n^{-2} shows a linear behavior for the data points for which the condition $J_n \geq J_{\text{lim}}$ is satisfied which confirms the Sand behavior. As an example, **Supplementary Figure 6** shows τ_s vs. J_n^{-2} for a LMP(48-18) battery including the linear regression line. The slope of the regression line is then directly proportional to D_{amb} and thus to D_{Li^+} . Considering a t^+ of 0.15 as determined by impedance spectroscopy, D_{Li^+} has been calculated for each LMP(x - y) battery. D_{Li^+} is independent of the positive electrode and SPE thickness with an average value of $3.1 \pm 0.6 \cdot 10^{-8} \text{ cm}^2 \cdot \text{s}^{-1}$. So, the Li^+ diffusion coefficient determined by the Sand equation when applied to battery cycling data is rather similar to the Li^+ diffusion coefficient within the SPE. Therefore, the power limiting phenomena in these all-solid-state batteries is then the diffusion of the Li^+ cation in the electrolyte spanning from the Li negative electrode to the Al current collector of the positive electrode rather than the diffusion of Li^+ in the LFP active material (Doyle and Newman, 1995). The use of the sand equation is then an efficient tool to quickly determine the effective limiting diffusion process at stake in batteries when $J > J_{\text{lim}}$, i.e., at rates where only a part of the full capacity is recovered.

To go further and compare all the LMP(x - y) batteries together, J_{lim} was extracted from each Q_n/Q_0 vs. J_n plots. **Figure 4** represents the average normalized capacity as a function of the ratio J_{lim}/J_n for the LMP(20-18), LMP(33-36), LMP(48-216), and LMP(60-54) batteries. All the experimental normalized capacity curves of the LMP(x - y) batteries overlap to a mere curve which demonstrates that the limiting phenomenon in stake here is the same whatever the electrode and electrolyte thickness and is due to the diffusion in the polymer electrolyte. In literature, a mere curve is usually obtained when the normalized capacity is represented as a function of the C-rate or its inverse. The mere curve is then fitted by an exponential decay function stretched by



an empirical value, typically of 2, and adjusted using a relaxation time parameter (Heubner et al., 2018b; Tian et al., 2019). In our case, no semi-empirical function is used as the only parameter corresponds to a physical parameter J_{lim} , graphically determined. The next step is to provide a simple model to get more insight on the partition of J_{lim} due to the diffusion into the electrolyte and/or into the electrolyte spanning into the positive electrode.

Figure 5 represents the average J_{lim} of the different LMP(x - y) batteries as a function of the positive electrode thickness x . For a given positive electrode thickness, J_{lim} increases with the diminution of the SPE thickness. The J_{lim} values span within those obtained for a 20 μm thick positive electrode between 0.06 and 0.48 $\text{mA}\cdot\text{cm}^{-2}$ when y equals to 216 and 18 μm , respectively. When $y = 18 \mu\text{m}$, J_{lim} decreases linearly with x . For higher y value, the slope of the J_{lim} decays with x is less pronounced as y increases until a plateau value for the highest reported SPE thicknesses of 216 μm . The best power performance is thus expected for battery comprising a thin positive electrode and a thin SPE layer. However, by looking at the interplay between the x and y value, battery with higher specific energy can be assembled without impairing much the power performances. Indeed, J_{lim} is an indicator of the onset of battery capacity decrease from its nominal value, i.e., Q_0 . In **Figure 5**, the LMP(60-18) battery shows a J_{lim} value higher than the one of the LMP(20-36). This implies that J_{lim} is a weak function of the positive electrode thickness and a strong function of the SPE thickness. In consequence, for the LMP(x - y) batteries, the specific volumetric energy density per single cell can be augmented simply by selecting the best compromise between positive electrode and SPE thicknesses.

The weak dependence of J_{lim} relatively to x can be understood by considering that the effective surface ensuring the Li^+ diffusive

flux at the interface of composite electrode/electrolyte is lower than the geometrical surface of the electrode due to the presence of active material and carbon particles in the composite electrode. So, the current density corresponding to the flux of Li^+ in the SPE located inside the electrode is higher than the one in the electrolyte. In addition, the Li^+ diffusion length depends on the electrode tortuosity. These effects can be described by introducing a parameter α with $0 < \alpha \leq 1$ in which the total Li^+ diffusion path corresponds to $\alpha \cdot x + y$. One can expect that the diffusion path in the composite electrode should be longer than the thickness of the electrode due to its tortuosity. However, here we measure an effective value, and it appears that the diffusion in the electrolyte located in the electrode seems higher than that of the raw electrolyte probably due to moisture content. This leads to an α value below unity. In **Figure 6**, J_{lim} is plotted as a function of $\alpha \cdot x + y$. By adjusting α to a value of 0.35, all the J_{lim} data collapse onto a master curve that is fitted using a reciprocal function, shown as the dashed curve in **Figure 6**, that yields to the following equation:

$$J_{lim} = \frac{K}{(\alpha \cdot x + y)} \tag{8}$$

with $K = 13.2 \text{ mA}\cdot\text{cm}^{-1}$ based on least squares fit ($R^2 > 0.99$).

To provide a physical meaning of the parameter K in Equation (8), the LMP(x - y) can be modeled via a one dimensional approach as extensively performed by the group of Doyle et al. (1993). Thus, the Fick law in 1-D can be simply applied in which the diffusion current is proportional to the Li^+ concentration gradient, the thickness of the diffusion process, and the diffusion coefficient of the limiting phenomena. Based on the results

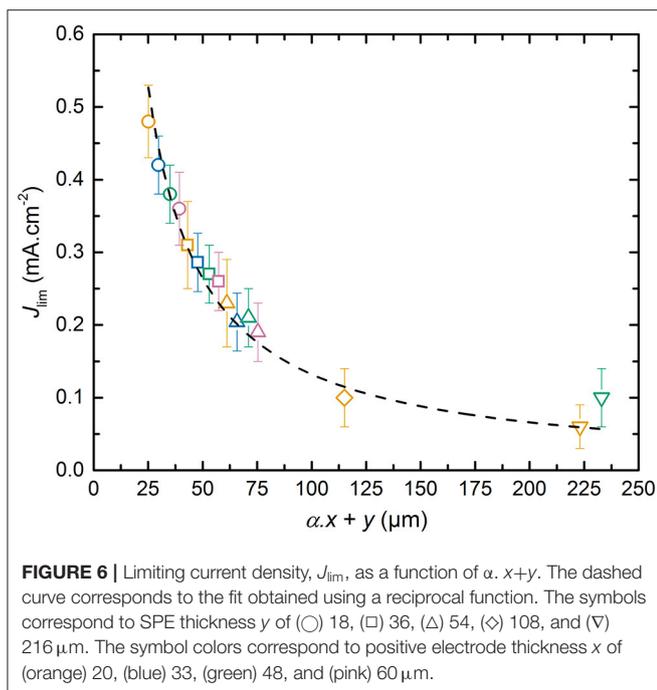


FIGURE 6 | Limiting current density, J_{lim} , as a function of $\alpha \cdot x + y$. The dashed curve corresponds to the fit obtained using a reciprocal function. The symbols correspond to SPE thickness y of (○) 18, (□) 36, (△) 54, (◇) 108, and (▽) 216 μm . The symbol colors correspond to positive electrode thickness x of (orange) 20, (blue) 33, (green) 48, and (pink) 60 μm .

obtained in **Figure 6**, J_{lim} is an inverse function of the sum $\alpha \cdot x + y$ and corresponds thus to the current value at which the Li concentration gradient throughout the battery falls down to a null value at the Al current collector/positive electrode. J_{lim} can then be expressed as a function of the effective Li^+ diffusion coefficient (D_{Li}^{eff}) according to :

$$J_{lim} = n \cdot F \cdot D_{Li}^{eff} \cdot \left(\frac{C_{Li}}{\alpha \cdot x + y} \right) \quad (9)$$

Combining Equation (9) into (8) permits to directly calculate a corresponding D_{Li}^{eff} value of $1.6 \cdot 10^{-8} \text{ cm}^2 \cdot \text{s}^{-1}$ based on the K parameter. The effective diffusion coefficient is then in the same range than the Li^+ diffusion coefficient determined by the Sand time methodology. In consequence, the effective diffusion coefficient D_{Li}^{eff} corresponds to the effective Li^+ diffusion coefficient D_{Li^+} in the SPE from the positive electrode to electrolyte layer. The J_{lim} methodology is less accurate than the Sand time methodology but is much faster to implement in order to quickly determine the limiting diffusion phenomena in the batteries. In addition, the α parameter is interesting since it depends on the electrode formulation and tortuosity as well as the ionic transport within the SPE used in the positive electrode. The variability of this parameters as function of the electrode formulation is a tool to optimize the electrode composition and texture, but is beyond the scope of our study.

CONCLUSION

The battery power signature is rapidly determined via a rapid power test. This method consists in applying consecutive

galvanostatic discharge step from high to low current density. The results given by this method is identical to that obtained by conventional galvanostatic cycling, i.e., series of charge-discharge steps. At current density higher than the limiting current density (J_{lim}), the application of the Sand equation for controlled-current technique is proven to be effective in determining the diffusion coefficient of the limiting process. Herein, in the Li polymer batteries, the Li^+ diffusion in the solid polymer electrolyte acting as battery separator and positive electrode binder limits the battery performance. In addition to the fast determination of J_{lim} , a physical meaning of this parameter is provided. J_{lim} is directly linked to an effective Li^+ diffusion coefficient throughout the battery that is similar to the one reported by the Sand equation. In consequence, the rapid power test is an efficient technique to compare series of batteries that differ by their assembly and to determine the main limiting factor, and thus optimize in turn the battery assembly. Battery design as then to consider the cumulative effect of both effect arising from the SPE layer thickness and from the positive electrode thickness. For example, in our LMP batteries, for a 54 μm thick electrolyte a positive electrode up to 48 μm can be used without impairing the battery power performance. For completeness, other factors than the power performance should also be considered, especially in case of Li metal based battery with the impact of the heterogeneous Li electrodeposition during charge step. To go further, this methodology should be applied to Li metal polymer batteries having different positive electrode formulations as well as on other battery technology.

DATA AVAILABILITY STATEMENT

All datasets generated for this study are included in the article/**Supplementary Material**.

AUTHOR CONTRIBUTIONS

RB, DD, MD, and ML designed the project study. MD and ML produced the composite electrodes. DD, PD, and HL conducted the experiments. RB, DD, and HL analyzed the data. RB and DD wrote the manuscript. All the authors commented the manuscript.

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SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/article/10.3389/fenrg.2019.00168/full#supplementary-material>

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Conflict of Interest: MD and ML are employed by the company Blue Solutions.

The remaining authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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